# CHAOTIC SOLUTION FOR THE BLACK-SCHOLES EQUATION 

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#### Abstract

The Black-Scholes semigroup is studied on spaces of continuous functions on $(0, \infty)$ which may grow at both 0 and at $\infty$, which is important since the standard initial value is an unbounded function. We prove that in the Banach spaces $$
Y^{s, \tau}:=\left\{u \in C((0, \infty)): \lim _{x \rightarrow \infty} \frac{u(x)}{1+x^{s}}=0, \lim _{x \rightarrow 0} \frac{u(x)}{1+x^{-\tau}}=0\right\}
$$ with norm $\|u\|_{Y^{s}, \tau}=\sup _{x>0}\left|\frac{u(x)}{\left(1+x^{s}\right)\left(1+x^{-\tau}\right)}\right|<\infty$, the Black-Scholes semigroup is strongly continuous and chaotic for $s>1, \tau \geq 0$ with $s \nu>1$, where $\sqrt{2} \nu$ is the volatility. The proof relies on the Godefroy-Shapiro hypercyclicity criterion.


## 1. Introduction

In B-S, F. Black and M. Scholes proved that under certain assumptions about the market, the value of a stock option, as a function of the current value of the underlying asset $x \in \mathbb{R}^{+}=[0,+\infty)$ and time, $u(x, t)$, satisfies the final value problem
(BS)

$$
\begin{cases}\frac{\partial u}{\partial t}=-\frac{1}{2} \sigma^{2} x^{2} \frac{\partial^{2} u}{\partial x^{2}}-r x \frac{\partial u}{\partial x}+r u & \text { in } \mathbb{R}^{+} \times[0, T] ; \\ u(0, t)=0 & \text { for } t \in[0, T] ; \\ u(x, T)=(x-p)^{+} & \text {for } x \in \mathbb{R}^{+},\end{cases}
$$

where $p>0$ represents a given strike price, $\sigma>0$ is the volatility and $r>0$ is the interest rate.

Let $v(x, t)=u(x, T-t)$. Then $v$ satisfies the forward Black-Scholes equation, which is a parabolic problem, defined for all time $t \in \mathbb{R}^{+}$by
(FBS)

$$
\begin{cases}\frac{\partial v}{\partial t}=\frac{1}{2} \sigma^{2} x^{2} \frac{\partial^{2} v}{\partial x^{2}}+r x \frac{\partial v}{\partial x}-r v & \text { in } \mathbb{R}^{+} \times \mathbb{R}^{+} \\ v(0, t)=0 & \text { for } t \in \mathbb{R}^{+} \\ v(x, 0)=f(x) & \text { for } x \in \mathbb{R}^{+}\end{cases}
$$

Strictly speaking, the condition $t \in \mathbb{R}^{+}$should have been written as $0 \leq t \leq T$. But once one notes that, there is no problem considering all nonnegative values of time.

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In (FBS) we have

$$
f(x)=(x-p)^{+}= \begin{cases}x-p & \text { if } x>p  \tag{1.1}\\ 0 & \text { if } x \leq p\end{cases}
$$

but for the time being we prefer to consider $f$ merely as an arbitrary given function. Later we shall deal with (1.1). In order to put the (FBS) problem in an abstract form, let us denote by $D_{\nu}=\nu x \frac{d}{d x}$, where $\nu=\sigma / \sqrt{2}$, and let

$$
\begin{equation*}
\mathcal{B}=D_{\nu}^{2}+\gamma D_{\nu}-r I=\nu^{2} C_{1}+C_{2}, \tag{1.2}
\end{equation*}
$$

with $\gamma=r / \nu-\nu, C_{1}:=x^{2} \frac{d^{2}}{d x^{2}}=D_{1}^{2}-D_{1}$ and $C_{2}:=r D_{1}-r I$. Then (FBS) can be written as
(AFBS)

$$
\left\{\begin{array}{l}
d v / d t=\mathcal{B} v \\
v(0, t)=0 \\
v(x, 0)=f(x) \quad \text { for } \quad x \in \mathbb{R}^{+}
\end{array}\right.
$$

For European call options, Cruz-Báez and González-Rodríguez C-G1 and Arendt and de Pagter AdP] showed that (FBS) is governed by a $C_{0}$-semigroup on a suitable Banach space. In [C-G2] the authors have generalized [C-G1 to American call options, a topic of interest in mathematical finance. But by working in the context of a contraction semigroup, these authors could not consider the issue of chaos. Recently, GMR gave a simple explicit representation of the solution of (FBS), and this representation holds in the spaces $Y^{s, \tau}$ considered here.

## 2. Multiplicative ( $C_{0}$ ) semigroups on the weighted space

For representing the Black-Scholes semigroup, we begin by introducing the translation on the multiplication group of positive numbers, $G=((0, \infty), \cdot)$. We do this now and we postpone the Black-Scholes semigroup to Section 3.

Let $\mu$ be the Haar measure on $G$ and suppose $\tau=\left\{\tau_{t}: t \in \mathbb{R}\right\}$ is the group of translations on $G$. Thus $d \mu=\frac{d x}{x}$, and $\tau_{t}(x)=\mathrm{e}^{t} x$, for $x>0, t \in \mathbb{R}$.

Let $s, \tau \geq 0$ and let $C(0, \infty)$ be the space of all complex continuous functions on $(0, \infty)$. Define

$$
Y^{s, \tau}:=\left\{u \in C(0, \infty): \lim _{x \rightarrow 0} \frac{u(x)}{1+x^{-\tau}}=\lim _{x \rightarrow \infty} \frac{u(x)}{1+x^{s}}=0\right\}
$$

with norm

$$
\|u\|_{s, \tau}=\sup _{x>0}\left|\frac{u(x)}{\left(1+x^{-\tau}\right)\left(1+x^{s}\right)}\right|<\infty
$$

These are Banach spaces.
Fix $\nu \in \mathbb{R} \backslash\{0\}$. Define the translation group with parameter $\nu, S_{\nu}:=\left\{S_{\nu}(t)\right.$ : $t \in \mathbb{R}\}$ on $Y^{s, \tau}$, by

$$
\left(S_{\nu}(t) f\right)(x)=\mathrm{e}^{t D_{\nu}} f(x)=f\left(\tau_{\nu t}(x)\right)
$$

for $f \in Y^{s, \tau}, x \in G$ and $t \in \mathbb{R}$. Since $\tau_{t+s}=\tau_{t} \tau_{s}$ for all $t, s \in \mathbb{R}, S_{\nu}$ forms a one-parameter group on each $Y^{s, \tau}$. Let $D_{\nu}$ be its infinitesimal operator in the sense of Hille, that is,

$$
D_{\nu} f=\left.\frac{d}{d t} S_{\nu}(t) f\right|_{t=0}
$$

for all $f$ for which this limit exists in $Y^{s, \tau}$; call this set $\mathcal{D}\left(D_{\nu}\right)$. Then $f \in \mathcal{D}\left(D_{\nu}\right)$ requires that $x \rightarrow f(x)$ and $x \rightarrow x f(x)$ are both in $Y^{s, \tau}$. Below we will establish
the strong continuity and characterize $\mathcal{D}\left(D_{\nu}\right)=\mathcal{D}\left(D_{1}\right)$ in $Y^{s, \tau}$ for $s \geq 1, \tau \geq 0$. The spaces $Y^{s, \tau}$ are Banach spaces, for all $s \geq 0, \tau \geq 0$.

Let $\mathcal{M}(0, \infty)$ be the set of all finite complex Borel measures on $(0, \infty)$. Any $\psi \in \mathcal{M}(0, \infty)$ can be written as

$$
\begin{equation*}
\psi=\operatorname{Re}(\psi)+\mathrm{i} \operatorname{Im}(\psi)=\sum_{j=1}^{4} c_{j} P_{j} \tag{2.1}
\end{equation*}
$$

where each $P_{j}$ is a probability measure on $(0, \infty)$ and the scalars $c_{j}$ satisfy $\operatorname{Re}(\psi)=$ $(\operatorname{Re} \psi)_{+}-(\operatorname{Re} \psi)_{-}=\psi_{1}-\psi_{2}$, with $\psi_{1}=c_{1} P_{1}$ and $\psi_{2}=-c_{2} P_{2}$, and $c_{1}, c_{2} \geq 0$. In the same way $\operatorname{Im}(\psi)=(\operatorname{Im} \psi)_{+}-(\operatorname{Im} \psi)_{-}=\psi_{3}-\psi_{4}$, with $\mathrm{i} \psi_{3}=c_{3} P_{3}$ and $\mathrm{i} \psi_{4}=-c_{4} P_{4}$, and $-\mathrm{i} c_{3},-\mathrm{i} c_{4} \geq 0$, and $P_{j}$ is uniquely determined for each $j$ for which $c_{j} \neq 0$. We also define $\psi \in \mathcal{M}_{l o c}(0, \infty)$ to mean that for any $n \in \mathbb{N}$, the restriction of $\psi$ to Borel subsets of $\left[\frac{1}{n}, n\right]$ is a finite complex Borel measure $\psi_{n}$ satisfying

$$
\begin{aligned}
\psi_{n} & =\operatorname{Re} \psi_{n}+\mathrm{i} \operatorname{Im} \psi_{n} \\
& =\left(\operatorname{Re} \psi_{n}\right)_{+}-\left(\operatorname{Re} \psi_{n}\right)_{-}+\mathrm{i}\left[\left(\operatorname{Im} \psi_{n}\right)_{+}-\left(\operatorname{Im} \psi_{n}\right)_{-}\right]
\end{aligned}
$$

Let $\xi_{n}$ denote any one of $\left(\operatorname{Re} \psi_{n}\right)_{ \pm},\left(\operatorname{Im} \psi_{n}\right)_{ \pm}$. Then each such $\xi_{n}$ determines uniquely a $\sigma$-finite Borel measure on $(0, \infty)$ via

$$
\xi(A)=\lim _{n \rightarrow \infty} \xi_{n}\left(A \cap\left[\frac{1}{n}, n\right]\right)
$$

for all Borel sets $A \subset[0, \infty]$. In this sense we can view $(\operatorname{Re} \psi)_{ \pm},(\operatorname{Im} \psi)_{ \pm}$as measures in a certain sense. Note that the set functions $\psi=(\operatorname{Re} \psi)_{+}-(\operatorname{Re} \psi)_{-}+\mathrm{i}\left[(\operatorname{Im} \psi)_{+}-\right.$ $\left.(\operatorname{Im} \psi)_{-}\right] \in \mathcal{M}_{l o c}(0, \infty)$ are not in general complex measures, but nevertheless we can treat them locally (away from 0 and $\infty$ ) as if they were complex measures by using $\psi_{n}$ for $n \in \mathbb{N}$.

We begin our study of $Y^{s, \tau}$ with the case of $s=0, \tau=0$. Note that

$$
Y^{0,0}=C_{0}(0, \infty)
$$

the continuous complex functions on $(0, \infty)$, which vanish at both 0 and $\infty$, with the norm

$$
\|u\|_{0,0}=\frac{1}{4}\|u\|_{\infty}
$$

for $u \in Y^{0,0}$. Note that the constant function 1 is in $Y^{s, \tau}$ if and only if $s>0, \tau>0$. By the Riesz Representation Theorem, the dual space of $Y^{0,0}=\left(\left(C_{0}(0, \infty),\|\cdot\|_{0,0}\right)\right.$ can be identified with $\mathcal{M}(0, \infty)$ with the norm

$$
\|\psi\|=4 T V(\psi)=4 \sum_{j=1}^{4}\left|c_{j}\right|
$$

when $c_{j}$ is as in (2.1), and $T V$ means total variation. The identification is made by mapping $u \in Y^{0,0}$ and $\psi \in \mathcal{M}(0, \infty)$ to

$$
\langle u, \psi\rangle=\int_{(0, \infty)} u(x) \psi(d x)
$$

We shall write $\int_{0}^{\infty}$ in place of $\int_{(0, \infty)}$.
One may view $Y^{0,0}$ as $\{u \in C[0, \infty]: u(0)=u(\infty)=0\}$, the continuous functions on the compact interval $[0, \infty]$, which vanish at both 0 and $\infty$. Similarly,
$\mathcal{M}(0, \infty)$ may be viewed as the finite complex Borel measures $\psi$ on $[0, \infty]$ satisfying $\psi(\{0, \infty\})=0$.

We next recall a well-known fact.
Lemma 2.1. Let $U: X \rightarrow Y$ be an isometric isomorphism between Banach spaces. Then $U^{*}: Y^{*} \rightarrow X^{*}$ is also an isometric isomorphism between their dual spaces.

Let

$$
\mathcal{C}_{c}:=\mathcal{C}_{c}(0, \infty)=\{u \in C(0, \infty): u \text { has compact support in }(0, \infty)\}
$$

Then $\mathcal{C}_{c}$ is dense in $Y^{s, \tau}$ for all $s, \tau \geq 0$. Let $u \in \mathcal{C}_{c}$. Then $u \in \mathcal{C}_{0}(\varepsilon, 1 / \varepsilon)$ for some $\varepsilon>0$. Let $\varphi \in \mathcal{M}([\varepsilon, 1 / \varepsilon])$ be a finite complex measure on $[\varepsilon, 1 / \varepsilon]$, which is the dual space of $\mathcal{C}[\varepsilon, 1 / \varepsilon]$. Then if $\psi \in\left(Y^{s, \tau}\right)^{*}$, we have $\langle u, \psi\rangle=\int_{0}^{\infty} u(x) \varphi(d x)$ for some $\varphi$ as above. We extend $\varphi$ by requiring that $\varphi(A)=0$ for all Borel subsets $A$ of $[0, \varepsilon] \cup[1 / \varepsilon, \infty]$. For this $\varphi$, define $\chi$ by

$$
\chi(d x)=\left(1+x^{s}\right)\left(1+x^{-\tau}\right) \varphi(d x)
$$

Then for $u \in Y^{s, \tau}, u=U_{s, \tau} v$ for a unique $v \in Y^{0,0}$, and

$$
\begin{aligned}
\langle u, \psi\rangle & :=\int_{0}^{\infty} u(x) \psi(d x)=\int_{0}^{\infty}\left(\frac{u(x)}{\left(1+x^{s}\right)\left(1+x^{-\tau}\right)}\right)\left(\left(1+x^{s}\right)\left(1+x^{-\tau}\right) \psi(d x)\right) \\
& =\left\langle 4 U_{s, \tau} u, \frac{1}{4} \chi\right\rangle=\left\langle U_{s, \tau} u, \chi\right\rangle=\langle v, \chi\rangle=\left\langle u, U_{s, \tau}^{*} \chi\right\rangle
\end{aligned}
$$

since $u \in \mathcal{C}_{c}$, which is dense in $Y^{a, b}$ for all $a, b>0$. Here $U_{s, \tau}$ is the $U$ of Lemma 2.1 corresponding to $X=Y_{s, \tau}$. Let
$Z_{s, \tau}:=\left\{\psi \in \mathcal{M}_{l o c}(0, \infty): \chi(d x):=\left(1+x^{s}\right)\left(1+x^{-\tau}\right) \psi(d x)\right.$ defines $\left.\chi \in \mathcal{M}(0, \infty)\right\}$
for $s, \tau \geq 0$. Then (2.2) below holds for $\psi \in Z_{s, \tau}$ and $\psi=U_{s, \tau}^{*} \chi$ for a unique $\chi \in \mathcal{M}(0, \infty)$, that is, $\psi \in Y_{s, \tau}^{*}$, and conversely. Thus we have proved that $Z_{s, \tau}$ can be identified with $\left(Y^{s, \tau}\right)^{*}$ for all $s, \tau \geq 0$. We restate this now proved result as follows.

Lemma 2.2. For $s, \tau \geq 0$, the dual space of $Y^{s, \tau}$ is (2.2)

$$
\left(Y^{s, \tau}\right)^{*}=\left\{\varphi \in \mathcal{M}_{l o c}((0, \infty)): \eta(d x):=\left(1+x^{s}\right)^{-1}\left(1+x^{-\tau}\right)^{-1} \varphi(d x) \in \mathcal{M}((0, \infty))\right\}
$$

Let us define the space $\mathscr{S}_{s, \tau}:=\left\{f \in C^{1}(0, \infty) \cap Y^{s, \tau}: f^{\prime} \in L^{\infty}(0, \infty)\right\}$. In order to prove that $S_{\nu}$ is a $\left(C_{0}\right)$ group on $Y^{s, \tau}$, we need the following lemma.

Lemma 2.3. The space $\mathscr{S}_{s, \tau}$ is dense in $Y^{s, \tau}$ for all $s, \tau \geq 0$.
Proof. Note that the map

$$
\frac{f(x)}{4} \longrightarrow \frac{f(x)}{\left(1+x^{s}\right)\left(1+x^{-\tau}\right)}
$$

is an isometric isomorphism from $Y^{0,0}$ onto $Y^{s, \tau}$ which leaves invariant $C_{c}^{\infty}(0, \infty)$, the smooth functions with compact support in $(0, \infty)$. Therefore $C_{c}^{\infty}(0, \infty)$ and $\mathscr{S}_{s, \tau}$ are both dense in $Y^{s, \tau}$ since $C_{c}^{\infty}(0, \infty)$ is dense in $Y^{0,0}$.

Theorem 2.4. The family $S_{\nu}$ forms a $\left(C_{0}\right)$ group on $Y^{s, \tau}$ for each $s \geq 1$ and $\tau \geq 0$.

Proof. First we note that the constant function 1 belongs to $Y^{s, \tau}$ if and only if $s$, $\tau$ are both positive. Next, we observe that for $f \in Y^{s, \tau}$ and $t \in \mathbb{R}$,

$$
\begin{aligned}
\left\|S_{\nu}(t) f\right\|_{s, \tau} & =\sup _{x>0} \frac{\left|f\left(\mathrm{e}^{\nu t} x\right)\right|}{\left(1+x^{s}\right)\left(1+x^{-\tau}\right)} \\
& =\sup _{y>0}\left|\frac{|f(y)|}{\left(1+\left[\mathrm{e}^{-\nu t} y\right]^{s}\right)\left(1+\left[\mathrm{e}^{-\nu t} y\right]^{-\tau}\right)}\right|
\end{aligned}
$$

Suppose $t \nu>0$. Then

$$
\left\|S_{\nu}(t) f\right\|_{s, \tau} \leq \mathrm{e}^{\nu t s} \sup _{y>0} \frac{|f(y)|}{\left(1+y^{s}\right)\left(1+y^{-\tau}\right)}=\mathrm{e}^{\nu t s}\|f\|_{s, \tau}
$$

For $t \nu \leq 0$, we have

$$
\begin{aligned}
\left\|S_{\nu}(t) f\right\|_{s, \tau} & =\sup _{y>0}\left|\frac{|f(y)|}{\left(1+\left[\mathrm{e}^{-\nu t} y\right]^{s}\right)\left(1+\left[\mathrm{e}^{-\nu t} y\right]^{-\tau}\right)}\right| \\
& \leq \mathrm{e}^{|\nu t| \tau} \sup _{y>0} \frac{|f(y)|}{\left(1+y^{s}\right)\left(1+y^{-\tau}\right)}=\mathrm{e}^{|\nu t| \tau}\|f\|_{s, \tau}
\end{aligned}
$$

Thus $S_{\nu}(t): Y^{s, \tau} \mapsto Y^{s, \tau}$ and $\left\|S_{\nu}(t)\right\| \leq \mathrm{e}^{\omega|t|}, \omega=|\nu| \max \{s, \tau\}$.
Thanks to Lemma 2.3, it is enough to show the strong continuity on $\mathscr{S}_{s, \tau}$. In fact, for any $f \in \mathscr{S}_{s, \tau}$, choose $\chi \in C^{\infty}(0, \infty)$ such that $\chi(x)=0$ for $0 \leq x \leq 1, \chi$ is increasing on $(1,2)$, and $\chi(x)=1$ for $x \geq 2$.

Let $f_{1}=f \chi, f_{2}=f(1-\chi)$. Then $f_{1}, f_{2} \in \mathscr{S}_{s, \tau}, \operatorname{supp} f_{1} \subset(1, \infty), \operatorname{supp} f_{2} \subset$ $(0,2)$, and $f_{1}+f_{2}=f$. Now for $f_{1}$,

$$
\begin{aligned}
\left\|S_{\nu}(t) f_{1}-f_{1}\right\|_{s, \tau} & =\sup _{x \geq 1} \frac{\left|f_{1}\left(\mathrm{e}^{\nu t} x\right)-f_{1}(x)\right|}{\left(1+x^{s}\right)\left(1+x^{-\tau}\right)} \\
& \leq\left\|f_{1}^{\prime}\right\|_{\infty} \sup _{x \geq 1} \frac{\left|\mathrm{e}^{\nu t} x-x\right|}{1+x^{s}} \\
& \leq\left\|f_{1}^{\prime}\right\|_{\infty}\left|\mathrm{e}^{\nu t}-1\right| \rightarrow 0, \quad \text { as } \quad t \rightarrow 0^{+} \text {since } s \geq 1
\end{aligned}
$$

For $f_{2}$, we have

$$
\left\|S_{\nu}(t) f_{2}-f_{2}\right\|_{s, \tau} \leq\left\|f_{2}^{\prime}\right\|_{\infty} \sup _{0<x<2} \frac{\left|\mathrm{e}^{\nu t} x-x\right|}{1+x^{-\tau}} \leq \frac{2^{\tau+1}}{1+2^{\tau}}\left\|f_{2}^{\prime}\right\|_{\infty}\left|\mathrm{e}^{\nu t}-1\right| \rightarrow 0
$$

as $t \rightarrow 0^{+}$, and this proves the theorem.
In the sequel we will need the following result, which is proved in G1 and deL Theorem 11].

Lemma 2.5. Suppose iA generates a strongly continuous group. Let $p(t)=t^{2 n}+$ $q(t)$, where $q$ is a polynomial of degree less than $2 n$. Then $-p(A)$ generates $a$ holomorphic ( $C_{0}$ ) semigroup of angle $\pi / 2$.

Take $A=-\mathrm{i} D_{\nu}$, so that $\mathrm{i} A$ generates a strongly continuous group on $X=Y^{s, \tau}$ and take $p(t)=t^{2}-\mathrm{i} \gamma t+r$. Hence we have the following result.

Theorem 2.6. The operator $\mathcal{B}$ defined in (1.2) generates a holomorphic $\left(C_{0}\right)$ semigroup of angle $\pi / 2$ on any $Y^{s, \tau}$, where $s \geq 1, \tau \geq 0$.

## 3. The chaotic character of the Black-Scholes semigroup

Let $X$ be a separable complex Banach space.
Definition 3.1. A strongly continuous semigroup (or $\left(C_{0}\right)$ semigroup) $T=\{T(t)$ : $t \geq 0\}$ of bounded linear operators on $X$ is called hypercyclic if there exists a vector $x \in X$ such that its orbit $\{T(t) x: t \geq 0\}$ is dense in $X$, and $T$ is called chaotic if in addition the set of periodic points of $T$,

$$
\mathcal{P}_{\text {per }}:=\left\{x \in X: \text { there exists } t_{0}>0 \text { such that } T\left(t_{0}\right) x=x\right\}
$$

is dense in $X$.
The notion of chaotic ( $C_{0}$ ) semigroups was introduced independently by MacCluer $[\mathrm{McC}]$ and Protopopescu and Azmy [P-A] ; the first systematic study of this concept is due to Desch, Schappacher and Webb DSW. So far, several specific examples of hypercyclic $\left(C_{0}\right)$ semigroups have come up in the literature (see GE1, GE2 for complete citations).

The following lemma is proved by G. Godefroy and J. Shapiro in G-S, Corollary 1.5$]$.

Lemma 3.2. Suppose $A$ is a linear bounded operator on a Banach space $X, Q_{1}, Q_{2}$ are dense subsets of $X$ and $Z: Q_{1} \mapsto Q_{1}$ such that
(1) $A Z y=y$, for all $y \in Q_{1}$,
(2) $\lim _{n \rightarrow \infty} Z^{n} y=0$, for all $y \in Q_{1}$ and
(3) $\lim _{n \rightarrow \infty} A^{n} w=0$, for all $w \in Q_{2}$.

Then $A$ is hypercyclic.
Let $s>1 / \nu$, where $\nu>0$ is given. Denote by

$$
\begin{equation*}
S_{s}=\{\lambda \in \mathbb{C}: 0<\operatorname{Re} \lambda<\nu s\} \tag{3.1}
\end{equation*}
$$

the open strip in $\mathbb{C}$ and let $h_{\lambda}(x)=x^{\lambda}$. This function is well-defined in $\mathbb{R}^{+}$for any $\lambda \in S_{s}$.

Lemma 3.3. The function $\lambda \mapsto h_{\lambda}(x)$ is analytic from $S_{s}$ into $Y^{s, \tau}$ for each $s \nu>1$ and $\tau \geq 0$.

Proof. Note that when $\tau=0$, any $\psi \in \mathcal{M}(0, \infty)$ cannot have an atom at 0 and $\mathbf{1} \notin Y^{s, 0}$. Now, since weak analyticity is equivalent to analyticity, we have only to prove that

$$
\lambda \mapsto \int_{(0, \infty)} h(x, \lambda) \psi(d x)
$$

is analytic for any $\psi \in F$, where $F$ is a norm-determining subset of $\left(Y^{s, \tau}\right)^{*}$. The norm-determining set we use is

$$
F:=\left\{c \delta_{x}: c \in \mathbb{C}, x \in(0, \infty)\right\}
$$

where $\delta_{x}$ denotes the Dirac point mass measure at $x$. Note that

$$
\|f\|_{s, \tau}=\sup \left\{|c f(x)|=|\langle f, \psi\rangle|: \psi=c \delta_{x}, c \in \mathbb{C}, x \in(0, \infty),\|\psi\|_{\left(Y^{s, \tau}\right)^{*}}=1\right\}
$$

and a choice of $c$ that works above is $c=\left[\left(1+x^{s}\right)\left(1+x^{-\tau}\right)\right]^{-1}$, when the supremum defining the norm of $f$ is a maximum attained at $x$. Furthermore

$$
\lambda \mapsto x^{\lambda}=\mathrm{e}^{(\ln x) \lambda}=\left\langle h_{\lambda}(x), \delta_{x}\right\rangle
$$

is an entire function of $\lambda \in \mathbb{C}$ for all $x>0$; and for $\lambda \in S_{s}$ and $x>0, \lim _{x \rightarrow \infty} x^{\lambda} /$ $\left(1+x^{s}\right)=0$. Hence $x^{\lambda} \in Y^{s, \tau}$ for all $\tau>0$.

If a linear operator $L$ generates a $\left(C_{0}\right)$ group on a Banach space $X$, then some polynomials in $L$ (such as $L^{2}+\alpha L+\beta I$ for arbitrary scalars $\alpha, \beta$ ) generate $\left(C_{0}\right)$ semigroups on $X$. For the operator $\mathcal{B}$, defined in (1.2), the Black-Scholes semigroup can be represented by $T(t):=f\left(D_{\nu}\right)$, where

$$
\begin{equation*}
f(z)=\mathrm{e}^{\operatorname{tg}(z)} \quad \text { with } g(z)=z^{2}+\gamma z-r \tag{3.2}
\end{equation*}
$$

According to Theorem [2.6 this $\left(C_{0}\right)$ semigroup is well defined for each $t \in \mathbb{C}$, with $\operatorname{Re}(t)>0$. These operators will be shown to be chaotic on $X=Y^{s, \tau}$ for $s>1, \tau \geq 0$ when $s \nu>1$. We begin by recalling the following lemma, which was proved in DSW] and dL-E], and we reproduce this proof in our case.

Lemma 3.4. Suppose that there exists a set $\Omega \subset S_{s}$ which has an accumulation point in $S_{s}$. Then

$$
Q:=\operatorname{Span}\left\{h_{\lambda} \quad: \quad \lambda \in \Omega\right\}
$$

is dense in $Y^{s, \tau}$ for $s>1, \tau \geq 0$.
Proof. Suppose $\psi \in Q^{\perp}$. Since $\psi$ belongs to the dual of $Y^{s, \tau}$ and $h_{\lambda}=x^{\lambda} \in Y^{s, \tau}$, Lemma 3.3 asserts that $p(\lambda)=\left\langle\psi, x^{\lambda}\right\rangle$ is well defined and $p(\lambda)$ is analytic in $S_{s}$. Since $p(\lambda)=0$ for all $\lambda \in \Omega$, which is a set with an accumulation point, then $p=0$ in all of $S_{s}$ and so $\psi=0$, as desired.

We continue to work in the spaces $Y^{s, \tau}, s>1, \tau \geq 0$ with $s \nu>1$.
Lemma 3.5. Let $\mathbb{D}$ be the unit disk in $\mathbb{C}$ and $\mathbb{T}$, the unit circle, be its boundary. The set $f\left(S_{s}\right) \cap \mathbb{T}$ is nonempty and possesses infinitely many accumulation points in the strip $S_{s}$, where $f$ is as in (3.2).

Proof. For $f(z)=\mathrm{e}^{t g(z)}$ with $t>0$, in order to have $f\left(S_{s}\right) \cap \mathbb{T} \neq \emptyset$ we must find $z \in S_{s}$ such that

$$
\operatorname{Re} g(z)=\operatorname{Re}\left(\nu^{2} z^{2}+\left(r-\nu^{2}\right) z-r\right)=\nu^{2}\left(x^{2}-y_{0}^{2}-x\right)+r x-r=0
$$

with $z=x+\mathrm{i} y_{0}$. Equivalently, we must find $\left(x, y_{0}\right)$ with $0<x<\nu s, y_{0} \in \mathbb{R}$ such that

$$
\begin{equation*}
x^{2}+\left(\frac{r}{\nu^{2}}-1\right) x-\frac{r}{\nu^{2}}=y_{0}^{2} \tag{3.3}
\end{equation*}
$$



Figure 1

Call $\mathscr{C}$ the curve represented by the graph of the quadratic function $y=x^{2}+$ $\left(\frac{r}{\nu^{2}}-1\right) x-\frac{r}{\nu^{2}}$. As Figure 1 shows, for $1<x<\nu s$, there are uncountably many points $(x, y)$ on the dashed portion of $\mathscr{C}$ with $y>0$. For each such point let $y_{0}=\sqrt{y}$. Then this gives uncountably many solutions of (3.3).

Now we can prove our main theorem.
Theorem 3.6. The Black-Scholes $\left(C_{0}\right)$ semigroup $T$ is chaotic in $Y^{s, \tau}$ for each $s>1, \tau \geq 0$ with $s \nu>1$.

Proof. First let us prove that the $\left(C_{0}\right)$ semigroup $T=\left\{T(t)=f\left(D_{\nu}\right)=\mathrm{e}^{t \mathcal{B}}: t \geq\right.$ $0\}$ is hypercyclic. For this we will use Lemma 3.2, taking

$$
\Omega_{1}=\left\{\lambda \in \frac{1}{\nu} S_{s}:|f(\nu \lambda)|>1\right\}, \quad \Omega_{2}=\left\{\lambda \in \frac{1}{\nu} S_{s}: \quad|f(\nu \lambda)|<1\right\}
$$

and

$$
Q_{j}:=\operatorname{Span}\left\{h_{\lambda} \quad: \quad \lambda \in \Omega_{j}\right\} \quad \text { for } j=1,2
$$

Now, let $z_{0} \in f\left(S_{s}\right) \cap \mathbb{T}$, since $f$ is holomorphic and nonconstant, $f\left(S_{s}\right)$ is an open set, and $\Omega_{1}=f\left(S_{s}\right) \cap\{z \in \mathbb{C}:|z|>1\}$ and $\Omega_{2}=f\left(S_{s}\right) \cap\{z \in \mathbb{C}:|z|<1\}$ are also open, and any point in $\Omega_{j}$ is an accumulation point. So according to Lemma 3.4 $Q_{j}$ is dense in $Y^{s, \tau}$ for $j=1,2$.

Let $A=f\left(D_{\nu}\right)$ and define $Z=\left(f\left(D_{\nu}\right)\right)^{-1}$ on $Q_{1}$ so that

$$
Z\left(\sum_{k=1}^{N} \alpha_{k} h_{\lambda_{k}}\right)=\sum_{k=1}^{N} \alpha_{k}\left(f\left(\nu \lambda_{k}\right)\right)^{-1} h_{\lambda_{k}}
$$

for $\lambda_{k} \in \Omega_{1}, \alpha_{k} \in \mathbb{C}$ and $N \in \mathbb{N}$. It is clear that for any $y=\sum_{k=1}^{N} \alpha_{k} h_{\lambda_{k}} \in Q_{1}$, we have $A Z y=y$. Furthermore for $\lambda_{k} \in \Omega_{1},\left|f\left(\nu \lambda_{k}\right)\right|>1$, and consequently

$$
\lim _{n \rightarrow \infty} Z^{n} y=\lim _{n \rightarrow \infty} \sum_{k=1}^{N} \alpha_{k}\left(f\left(\nu \lambda_{k}\right)\right)^{-n} h_{\lambda_{k}}=0
$$

Finally, for $w=\sum_{k=1}^{N} \alpha_{k} h_{\lambda_{k}} \in Q_{2}$ with $\left|f\left(\nu \lambda_{k}\right)\right|<1$ for each $k$,

$$
\lim _{n \rightarrow \infty} A^{n} y=\lim _{n \rightarrow \infty} \sum_{k=1}^{N} \alpha_{k} f\left(\nu \lambda_{k}\right)^{n} h_{\lambda_{k}}(x)=0
$$

These imply that the hypotheses of the Godefroy-Shapiro Lemma 3.2 are satisfied and $A$ is hypercyclic.

To see that $T(t)=f\left(D_{\nu}\right)$ is chaotic, we define $\Omega_{3}=\left\{\lambda \in \frac{1}{\nu} S_{s} \quad: \quad f(\nu \lambda) \in \mathrm{e}^{2 \pi \mathrm{i} \mathbb{Q}}\right\}$ and $Q_{3}:=\operatorname{Span}\left\{h_{\lambda}: \lambda \in \Omega_{3}\right\} . Q_{3}$ is contained in the set of all periodic points of $A=f\left(D_{\nu}\right)$. Suppose $f\left(\nu \lambda_{k}\right)=\mathrm{e}^{2 \pi \mathrm{i} n_{k} / m_{k}}$. Then for $y=\sum_{k=1}^{N} \alpha_{k} h_{\lambda_{k}}$ and $m=\prod_{k=1}^{N} m_{k}$, one has $f\left(D_{\nu}\right)^{m} y=y$. So the set of all periodic points $\mathcal{P}_{\text {per }}$ of $f\left(D_{\nu}\right)$ is dense, and consequently $T(t)$ is chaotic.

The real-world applications of (FBS) require nonnegative initial data and nonnegative solutions. The above proof that the Black-Scholes semigroup $T$ is chaotic uses holomorphic functions and thus requires the use of spaces of complex-valued functions. Theorem 3.6 would be more satisfying from an applied standpoint if it were valid for real functions. This is precisely the content of the next result.

Let $Y_{\mathbb{R}}^{s, \tau}$ be the real functions in $Y^{s, \tau}$. This is a real Banach space. If $f \in Y^{s, \tau}$, then by [GMR, eq. (17)], the solution of (FBS) is given by

$$
v(x, t)=(T(t) f)(x)=(4 \pi t)^{-1 / 2} \int_{-\infty}^{\infty} \mathrm{e}^{-y^{2} /(4 t)} f\left(x \mathrm{e}^{\left(r-\sigma^{2} / 2\right) t-(\sigma / \sqrt{2}) y}\right) d y
$$

Thus $T(t) f$ is real (resp., nonnegative) for each $t \geq 0$ if and only if $f$ is real (resp. nonnegative). Let $S_{T}$ be the restriction of $T$ to $Y_{\mathbb{R}}^{s, \tau}$. Then $S_{T}=\left\{S_{T}(t): t \geq 0\right\}$ is a $\left(C_{0}\right)$ semigroup on $Y_{\mathbb{R}}^{s, \tau}$ for $s \geq 1, \tau \geq 0$.

Theorem 3.7. The semigroup $S_{T}$ on $Y_{\mathbb{R}}^{s, \tau}$ is chaotic if $s>1$ and $\tau \geq 0$ when $s \nu>1$.

Proof. Let $f \in Y^{s, \tau}$ be given, where $s>1$ with $s \nu>1$, and $\tau \geq 0$. Let $g \in Y^{s, \tau}$ have a dense $T$-orbit. Then there is a sequence of times $t_{n} \rightarrow \infty$ such that $\left\|T\left(t_{n}\right) g-f\right\|_{s, \tau} \rightarrow 0$ as $n \rightarrow \infty$. Consequently, since $\operatorname{Re}(T(t) h)=T(t)(\operatorname{Re}(h))$ for all $h \in Y^{s, \tau}$,

$$
\begin{aligned}
\left.\| S_{T}\left(t_{n}\right)(\operatorname{Re}(g))-f\right) \|_{s, \tau} & \leq\left\|\sqrt{\left[S_{T}\left(t_{n}\right)(\operatorname{Re}(g))-f\right]^{2}+\left[S_{T}\left(t_{n}\right)(\operatorname{Im}(g))\right]^{2}}\right\|_{s, \tau} \\
& =\|\left[\operatorname{Re}\left(T\left(t_{n}\right) g-f\right]+\mathrm{i}\left[\operatorname{Im}\left(T\left(t_{n}\right) g\right)\right] \|_{s, \tau}\right. \\
& =\left\|T\left(t_{n}\right) g-f\right\|_{s, \tau} \rightarrow 0
\end{aligned}
$$

as $n \rightarrow \infty$. It follows that $S_{T}$ is hypercyclic.
Next, if $f$ is periodic of period $p$, then so are $\operatorname{Re}(f)$ and $\operatorname{Im}(f)$. Thus $S_{T}$ has a dense set of periodic points since $T$ does. The theorem follows.

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